# Pancharatnam-Zak phase for two-dimensional systems: Gauge-independent topological invariant Mohamadi, Sepide; Abouie, Jahanfar

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### Introduction

Topological phases are a wide range of quantum phases that are classified using quantized topological invariants. In 2D systems, the Chern number is used to find the topological phases, that is independent of gauge [1]. But in some systems that have time reversal and inversion symmetry, the Chern number is zero. In these systems, 2D Zak phase is non-zero [2]. But the Zak phase is gauge dependent, therefore defined 2D we Pancharatnam-Zak phase to detect the topological phases of 2D systems [3]. This invariant, like 1D Pancharatnam-Zak phase introduced earlier [3], is independent of gauge.

## Method

We first consider a two-dimensional system of electrons that are moving in a two-dimensional periodic potential  $V(\vec{r})$  and are under the influence of a weak electric field E:

$$\widehat{H}_{\overrightarrow{\alpha}}(t) = \left( \left( \overrightarrow{p} - \hbar \overrightarrow{\alpha}(t) \right)^2 / 2m \right) + V(\overrightarrow{r}),$$

Where  $\vec{p}$  and m are the momentum and mass of the electrons, respectively, and  $\vec{\alpha}(t) = -\vec{E}t/\hbar$ . The eigenstates of this Hamiltonian with the help of the Bloch's eigenstates can be wrote in the form  $|\Psi_{n\vec{k}_{m}\vec{\alpha}}\rangle = e^{i\vec{k}_{m}\cdot\vec{r}} |u_{n}(\vec{k}_{m}+\vec{\alpha})\rangle$ . Here, setting km = 0 without loss of generality. Therefore, the geometric phase obtained during the evolution from the initial state at time t<sub>i</sub> to the final state at time t<sub>f</sub> is as follows:

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$$\gamma_{g} = \sum_{n=1}^{occ.} \left[ \operatorname{Arg} \left\{ u_{n} \left( \vec{\alpha}(t_{i}) \right) + \left( i \int_{\vec{\alpha}(t_{i})}^{\vec{\alpha}(t_{i})} d\vec{\alpha}(t_{i}) \right) \right\} \right]$$

2D γ<sub>g</sub> is independent of gauge; If  $\gamma_g = 0$ , the system is in the trivial phase, and if  $\gamma_g = 2\pi$ , the system is in the non-trivial phase.

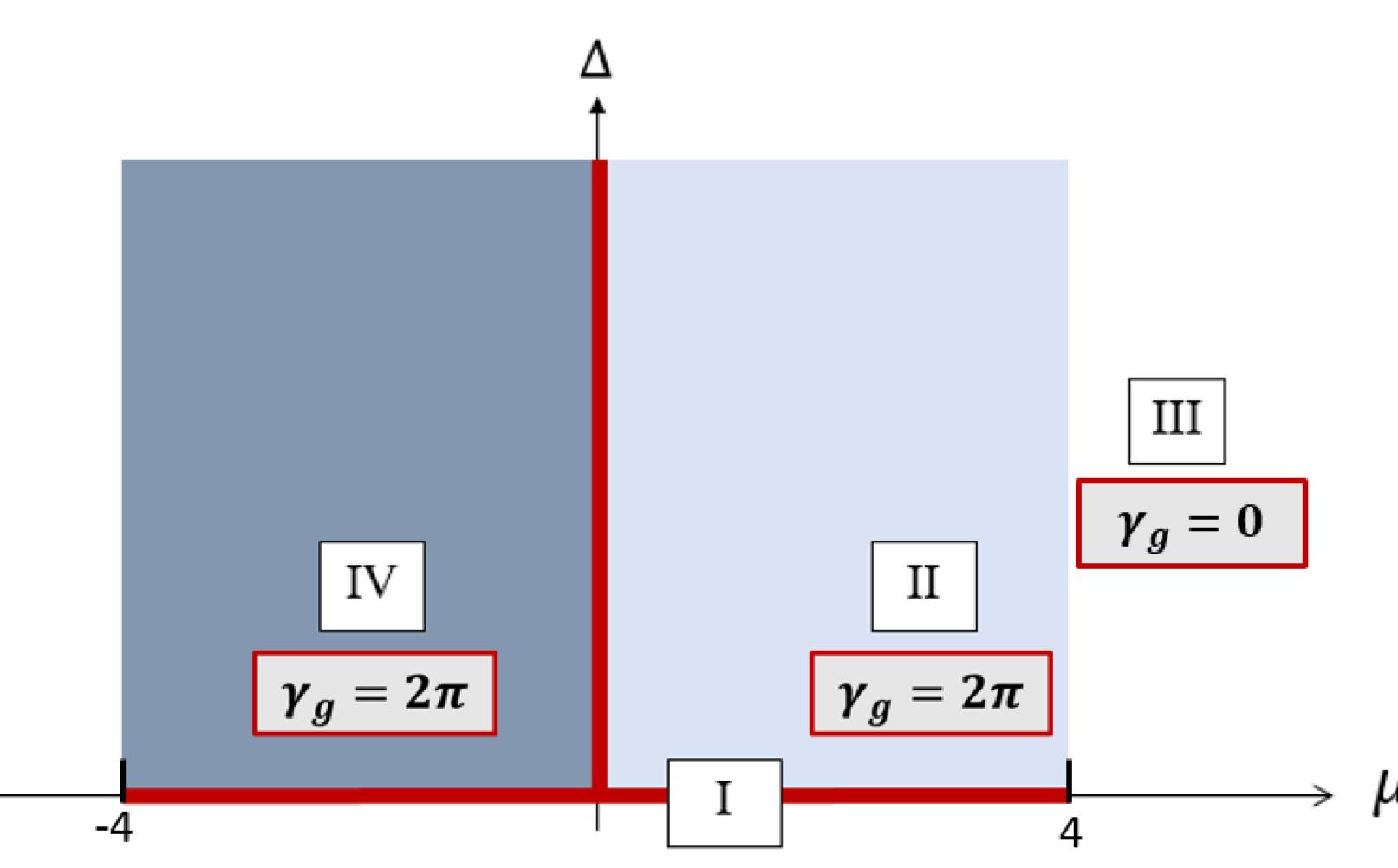


Figure 1: Phase diagram of 2D Kitaev model. According to this phase diagram, this model has four distinct phases:

Phase I: System is metal.

Phase II: System has topological non-trivial phase. Phase III: System is always gapped and is an insulator; and also the system has topological trivial phase. Phase IV : System has topological non-trivial phase.



 $\left| u_{n}\left( \vec{\alpha}(t_{f}) \right) \right|$  $|\vec{\alpha} \cdot \langle u_n(\vec{\alpha}) | \nabla_{\vec{\alpha}} | u_n(\vec{\alpha}) \rangle \Big|$ Pancharatnam-Zak phase, which is

### Results

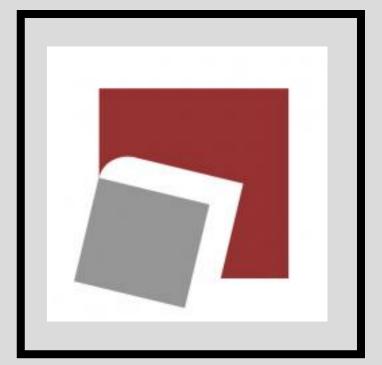
lf in 2D

 $\hat{H}_{Kitaev} = -J$ 

four distinct phases.

#### References





Pancharatnam-Zak phase, we use $|u_{n}(\vec{\alpha})\rangle e^{i\Lambda(\vec{\alpha})}$  instead of  $|u_{n}(\vec{\alpha})\rangle$ ,  $\gamma_{g}$  remains invariant, so  $\gamma_g$  is independent of gauge.

The 2D Kitaev model is used to describe the twodimensional p-wave superconductor. In this model J is the nearest neighbor hopping amplitude, and  $\mu$ represents the chemical potential and  $\Delta$  stands for the superconducting pairing amplitude:

$$\sum_{\vec{r},\vec{a}} \hat{c}_{\vec{r}}^{\dagger} \hat{c}_{\vec{r}+\vec{a}}^{\dagger} + h.c + \Delta \sum_{\vec{r},\vec{a}} \hat{c}_{\vec{r}} \hat{c}_{\vec{r}+\vec{a}}^{\dagger} + h.c$$
$$+ \mu \sum_{\vec{r},\vec{a}} \left( \hat{c}_{\vec{r}}^{\dagger} \hat{c}_{\vec{r}}^{\dagger} - (1/2) \right)$$

As an example, we have obtained the phase diagram of the 2D Kitaev model by examining the 2D Pancharatnam-Zak phase behavior in Figure 1.

According to this figure, the 2D Kitaev model has

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[3] N. Mukunda and R. Simon, "Quantum kinematic approach to the geometric phase. i. general formalism," Annals of Physics, vol. 228, no. 2, pp. 205–268, 1993.