

Rheology of Fluids: Newtonian to Non Newtonian

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Agenda:

- *Fluid: Definition*
- *Rheology: Elementary concepts*
- *Navier-Stokes Equation*
- *Basic Solutions*
- *Dilute Suspension*
- *Non Newtonian*
- *Oldroyd Model*

Fluid:

Deborah number, a measure of fluidity

$$De = \frac{\tau_{relax}}{\tau_{exp}}$$

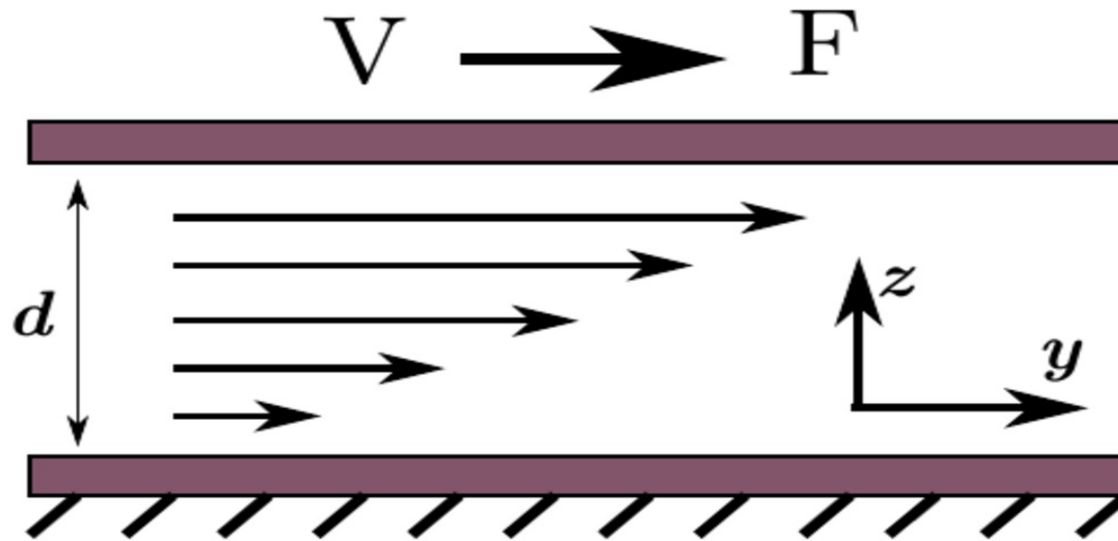
Relaxation time in response
to an externally applied force

Observation time



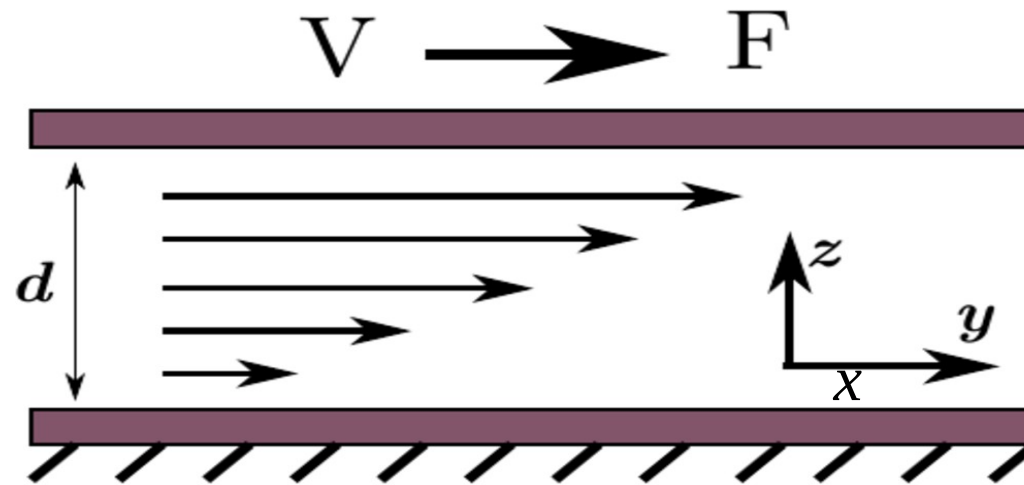
Rheology:

How does the system respond to an external force?



$$F = \Psi(V)$$

Viscosity: How does the system respond to an externally applied force?



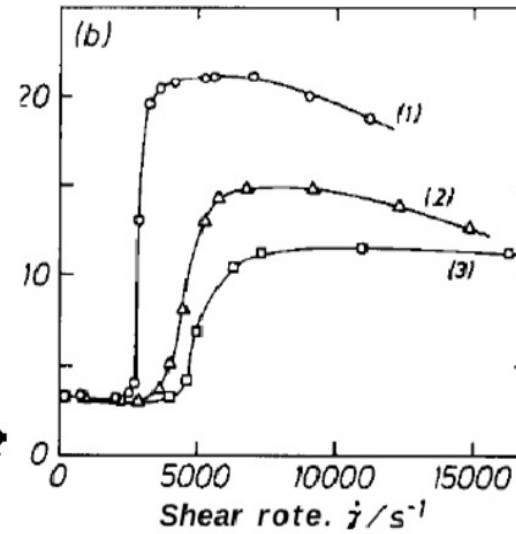
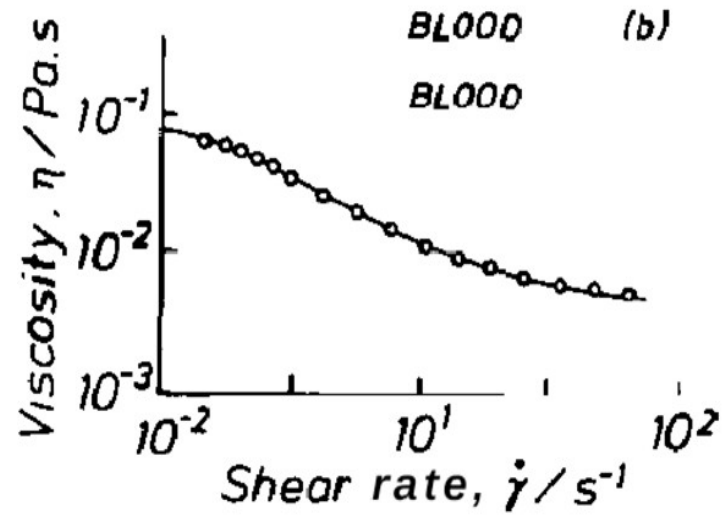
$$\sigma_{yz} := \frac{F}{A}$$

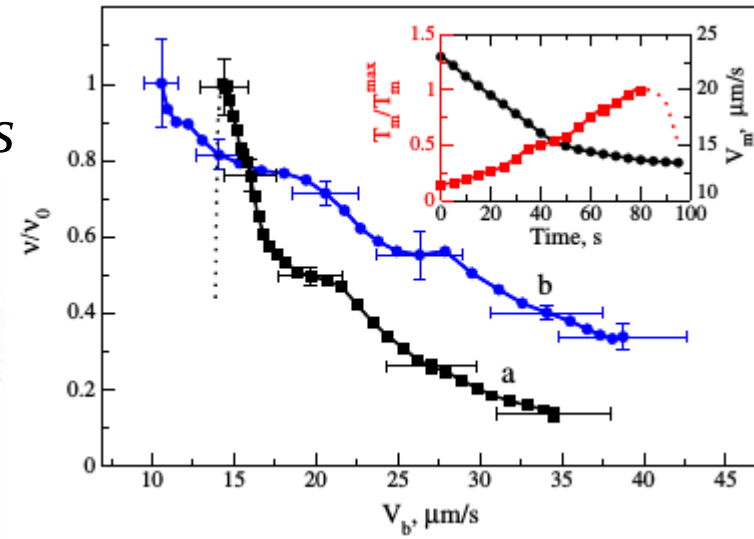
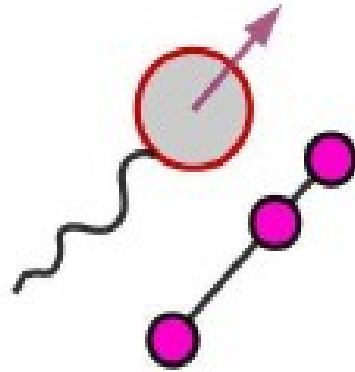
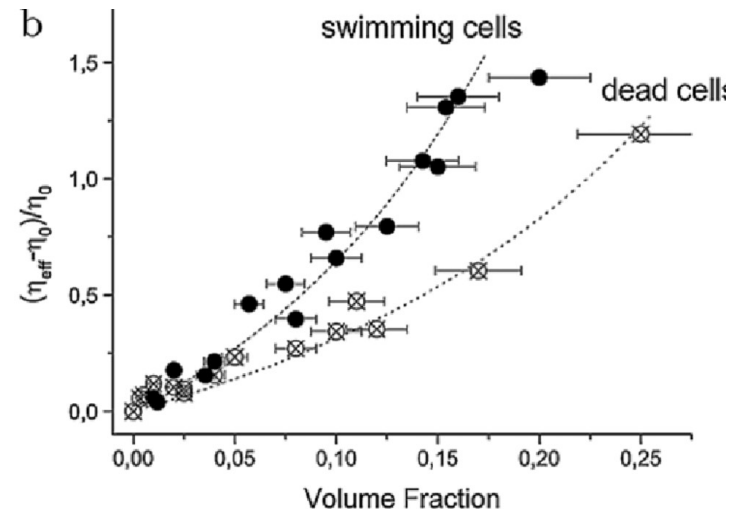
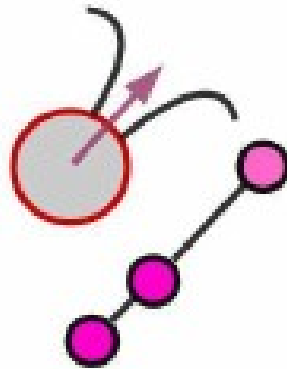
$$\sigma_{yz} = \Psi\left(\frac{\partial u_y}{\partial z}\right)$$

$$\frac{\partial u_y}{\partial z} = \dot{\gamma} = \frac{V - 0}{d}$$

Linear .and. isotropic → Newtonian
 Nonlinear .or. anisotropic → Non-Newtonian

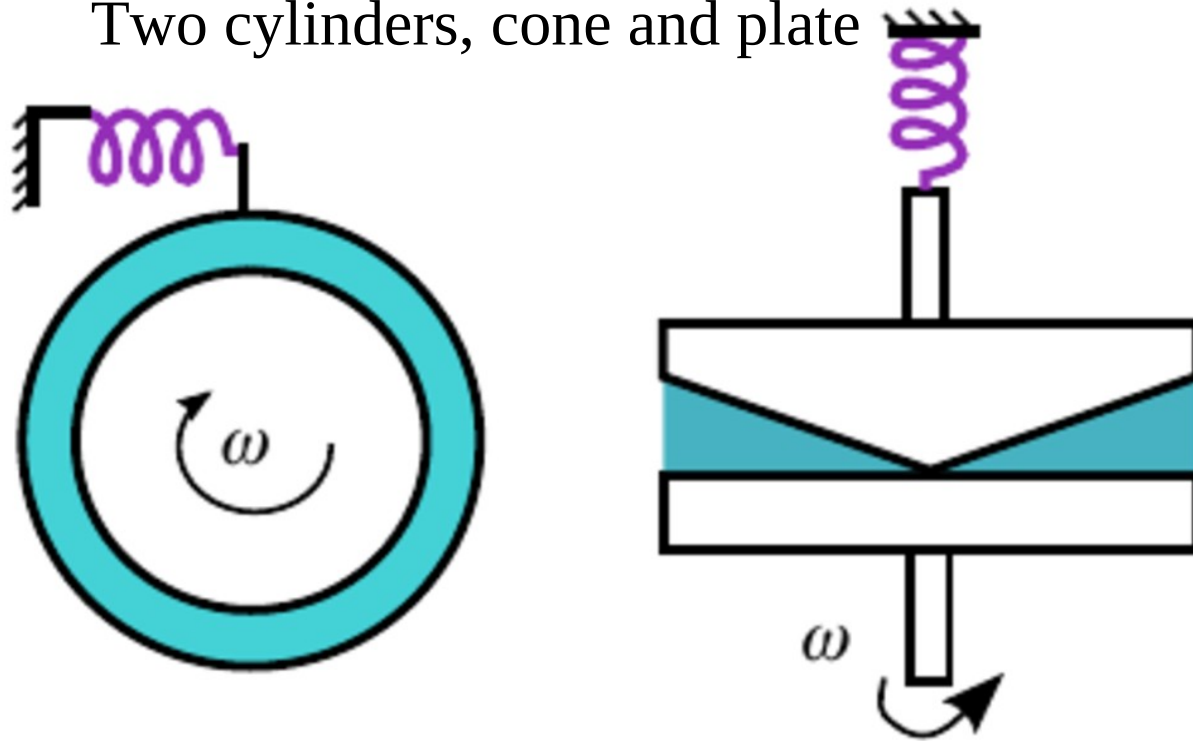
$$\sigma_{yz} = \eta \frac{\partial u_y}{\partial z} = \eta \dot{\gamma}$$

Some examples

Active Suspensions:*Bacillus Subtilis**Aranson, PRL 2009**Chlamydomonas**Rafai, PRL 2010*

Rheometry: Elementary tools

Two cylinders, cone and plate

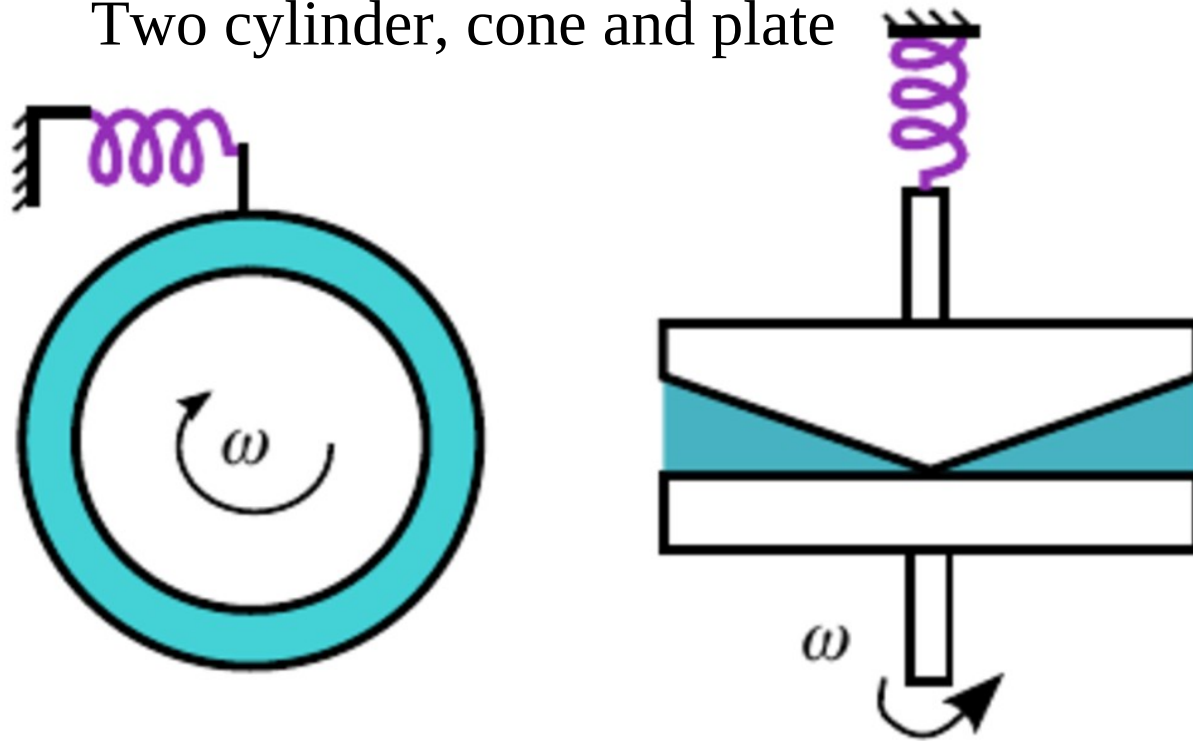


Dimensional Analysis for water

$$\eta \sim \frac{m}{lT}$$

Rheometry: Elementary tools

Two cylinder, cone and plate



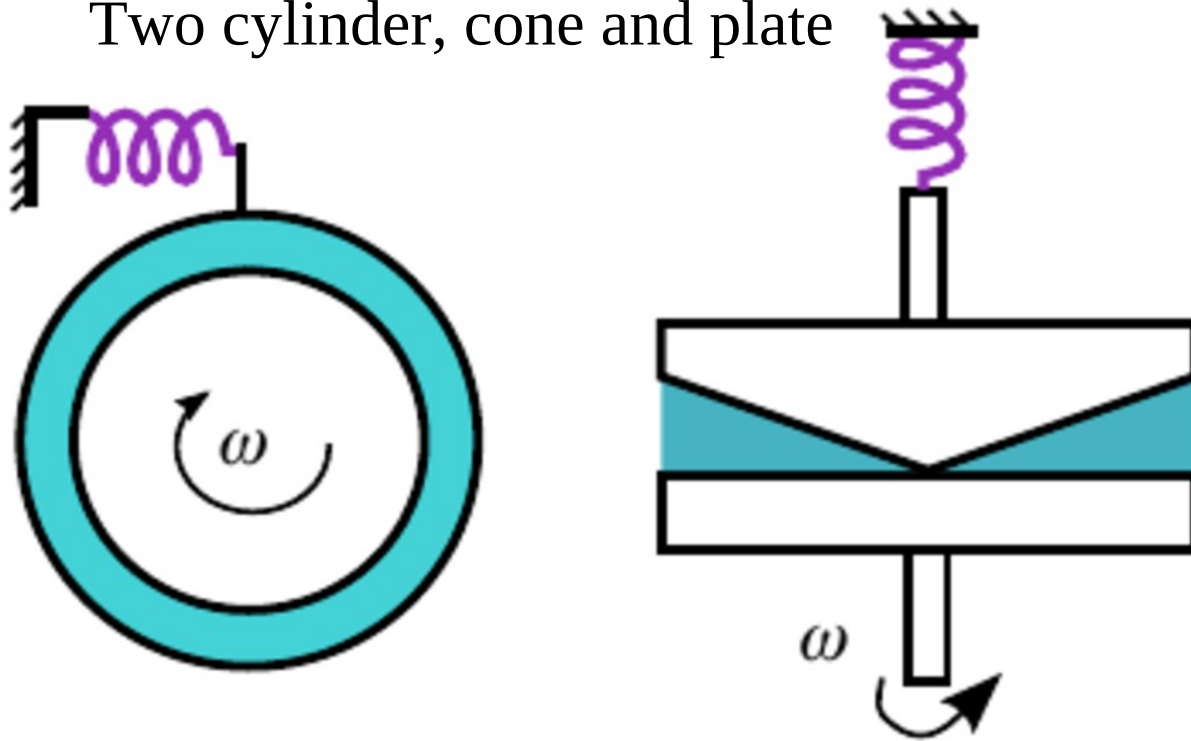
Dimensional Analysis for water

$$\eta \sim \frac{m}{lT} \rightarrow 10^{-26} \text{ kg}$$

$$10^{-9} \text{ m}$$

Rheometry: Elementary tools

Two cylinder, cone and plate



Dimensional Analysis for water

$$\eta \sim \frac{m}{lT} \rightarrow 10^{-29} \text{ kg}$$

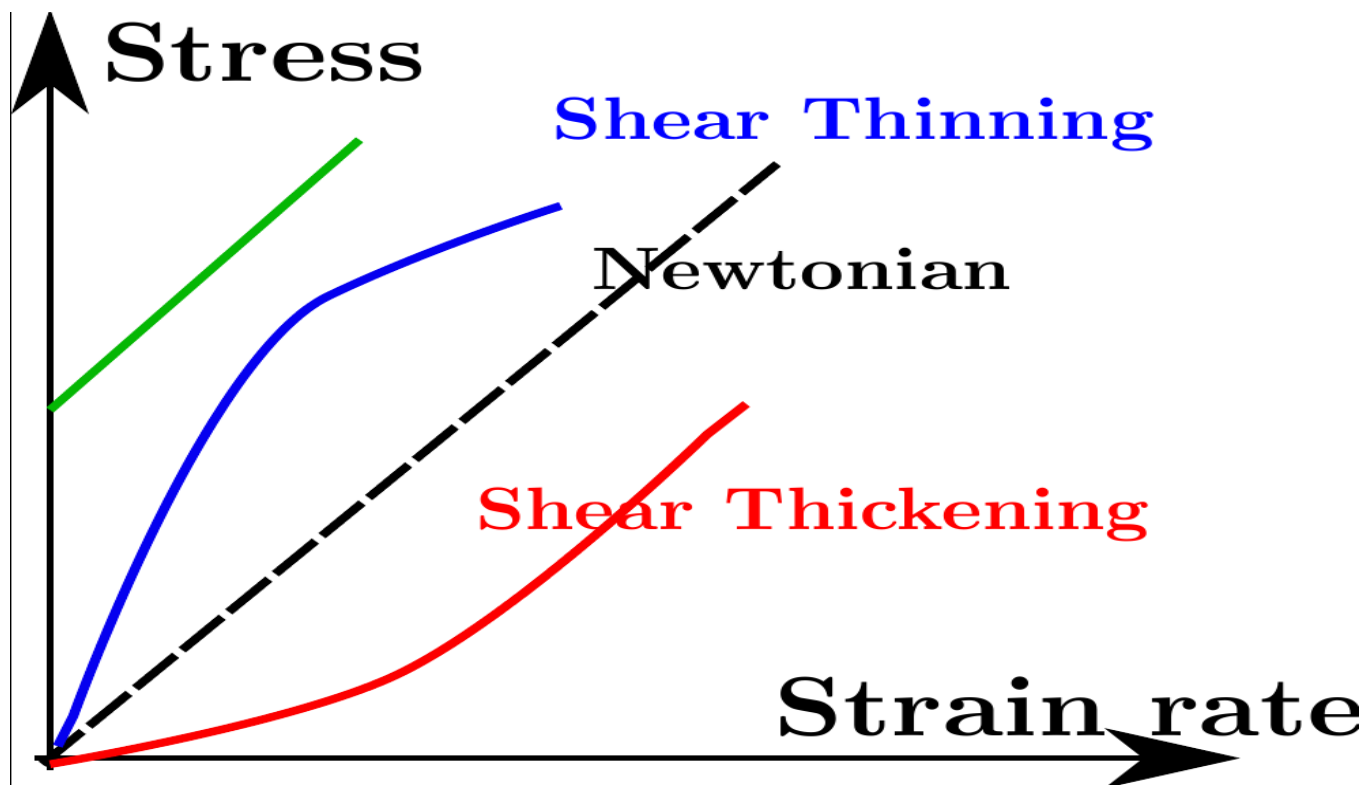
$$10^{-9} \text{ m}$$

$$\lambda \approx 1 \mu\text{m} \rightarrow \omega = c/\lambda \approx 10^8/10^{-6} = 10^{14}$$

$$T \approx 1/\omega = 10^{-14} \text{ s}$$

$$\eta \approx \frac{m}{lT} \approx 10^{-3} \text{ Pa}\cdot\text{s}$$

Non-Newtonian



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Toward the governing equations:

Velocity field $\mathbf{u}(\mathbf{x}, t)$

Acceleration field

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Velocity field $\mathbf{u}(\mathbf{x}, t)$

Acceleration field

$$a(\mathbf{x}, t) = (u(\mathbf{x} + \delta \mathbf{x}, t + \delta t) - u(\mathbf{x}, t)) / \delta t = \dot{\mathbf{u}}$$

$$\frac{d}{dt} \mathbf{u}(\mathbf{x}, t) = \left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) \mathbf{u}(\mathbf{x}, t)$$

Toward the governing equations:

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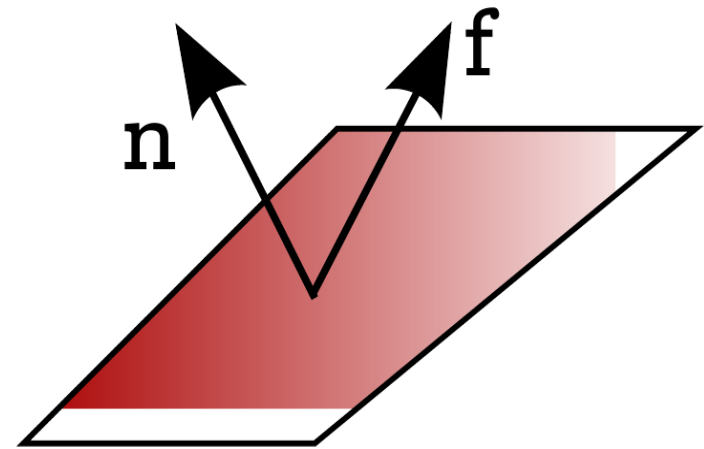
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$\sigma_{ij}(\mathbf{x}, t)$ Stress Tensor: i 'th component of the force exerted on a surface that is pointed to j direction

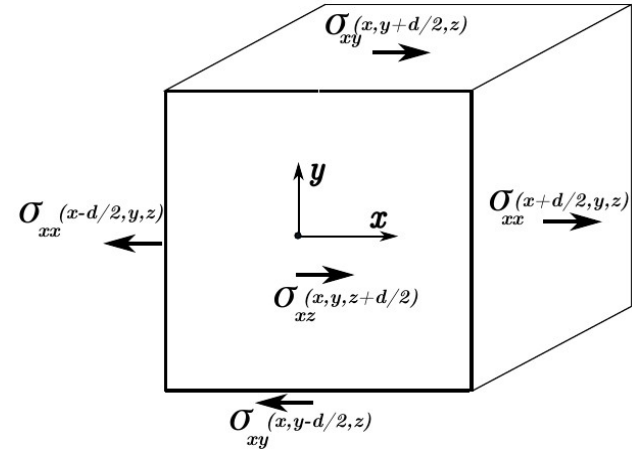
$$df_x = \sigma_{xx} ds_x + \sigma_{xy} ds_y + \sigma_{xz} ds_z$$

$$\mathbf{f} = \oint \boldsymbol{\sigma} \cdot d\mathbf{s}^T,$$



Governing equations

$$\mathbf{u}(\mathbf{x}, t) \quad \sigma_{ij}(\mathbf{x}, t)$$



$$\nabla \cdot (\rho \mathbf{u}) + \frac{\partial}{\partial t} \rho = 0 \quad \text{Continuity equation}$$

$$\rho(\mathbf{x}, t) \times dv \times \frac{d}{dt} \mathbf{u}(\mathbf{x}, t) = \mathbf{F}(\mathbf{x}, t) \times dv$$

$$F_x = d^3 \times \left[\sigma_{xx}\left(x + \frac{d}{\nu}, y, z\right) - \sigma_{xx}\left(x - \frac{d}{\nu}, y, z\right) + \sigma_{xy}\left(x, y + \frac{d}{\nu}, z\right) - \sigma_{xy}\left(x, y - \frac{d}{\nu}, z\right) + \sigma_{xz}\left(x, y, z + \frac{d}{\nu}\right) - \sigma_{xz}\left(x, y, z - \frac{d}{\nu}\right) \right]. \quad (4.4)$$

Governing Equations:

$$\sigma_{yz} = \eta \frac{\partial u_y}{\partial z} = \eta \dot{\gamma}$$

Phenomenological model for stress tensor

$$\sigma_{ij} = -P\delta_{ij} + \eta \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

$$\frac{d}{dt} \mathbf{u}(\mathbf{x}, t) = \left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) \mathbf{u}(\mathbf{x}, t)$$

Hydrodynamics:

Navier-Stokes Equation

$$\rho \left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) \mathbf{u}(\mathbf{x}, t) = \eta \nabla^2 \mathbf{u}(\mathbf{x}, t) - \nabla P(\mathbf{x}, t), \quad \nabla \cdot \mathbf{u} = 0$$

The Millennium Prize: 10^6 \$, The Clay Mathematics Institute of Cambridge, Massachusetts (CMI)

Colloid's Universe: Low Reynolds Regime:

$$\rho \left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) \mathbf{u} = -\nabla p + \eta \nabla^2 \mathbf{u}, \quad \nabla \cdot \mathbf{u} = 0$$

Colloid's Universe: Low Reynolds Regime:

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$$\mathcal{R}e = \frac{[\rho(\mathbf{u} \cdot \nabla)\mathbf{u}]}{\eta \nabla^2 \mathbf{u}} \approx \frac{\rho V/L}{\eta/L^2} = \frac{\rho V L}{\eta}$$

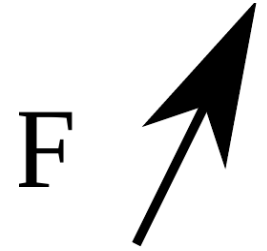
$$L = 0.1 \mu m, \quad V = 1 \mu m/S, \quad \eta = 10^{-3} PaS, \quad \rho = 10^3 Kg/m^3$$

$$Re \ll 1$$

$$\eta \nabla^2 \mathbf{u}(\mathbf{x}) - \nabla P(\mathbf{x}) = \mathbf{f}(\mathbf{x}), \quad \nabla \cdot \mathbf{u}(\mathbf{x}) = 0$$

Hydrodynamics:

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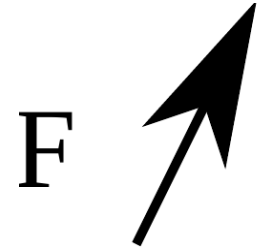
Point force:

$$\mathbf{f}(\mathbf{x}) = \mathbf{F} \delta(\mathbf{x})$$

$$u_i^\circ(\mathbf{x}) = \sum_j G_{ij} F_j, \quad G_{ij} = \frac{1}{4\pi\eta r} (\delta_{ij} + \hat{r}_i \hat{r}_j)$$

Hydrodynamics:

$$\eta \nabla^2 \mathbf{u}(\mathbf{x}) - \nabla P(\mathbf{x}) = \mathbf{f}(\mathbf{x}), \quad \nabla \cdot \mathbf{u}(\mathbf{x}) = 0$$

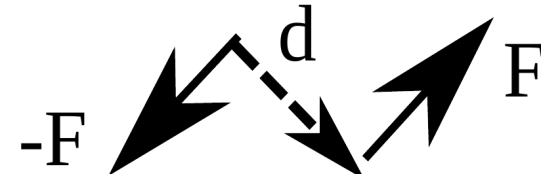
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Force Dipole:

$$\mathbf{u}^D = (\mathbf{G}(\mathbf{x} + \mathbf{d}) - \mathbf{G}(\mathbf{x})) \cdot \mathbf{F}$$



$$u_i^D = \sum_{j,k} G_{ijk}^D P_{kj} \quad P_{kj} = d_k F_j$$

$$G_{ijk}^D = \frac{\partial}{\partial x_k} G_{ij} = \frac{1}{4\pi\eta r^2} (\delta_{ij} \hat{r}_k - \delta_{ik} \hat{r}_j + \delta_{jk} \hat{r}_i + 3\hat{r}_i \hat{r}_j \hat{r}_k)$$

Multipole Expansion:

$$u_i(\mathbf{x}) = \sum_j G_{ij} b_j + \sum_{j,k} G_{ijk}^D p_{jk} + \sum_{j,k,l} G_{ijkl}^Q t_{jkl} + \text{چند قطبی‌های مراتب بالاتر}$$

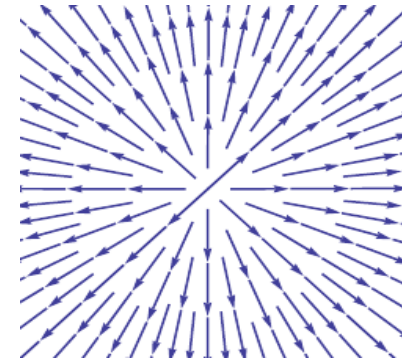
$$G_{ij} = \frac{1}{4\pi\eta r} (\delta_{ij} + \hat{r}_i \hat{r}_j)$$

$$G_{ijk}^D = \frac{\partial}{\partial x_k} G_{ij}$$

$$G_{ijlm}^Q = \frac{\partial}{\partial x_m} G_{ijl}^D$$

Sink and Source another singular solutions:

$$u_i^s = \frac{\hat{r}_i}{4\pi r^3}$$



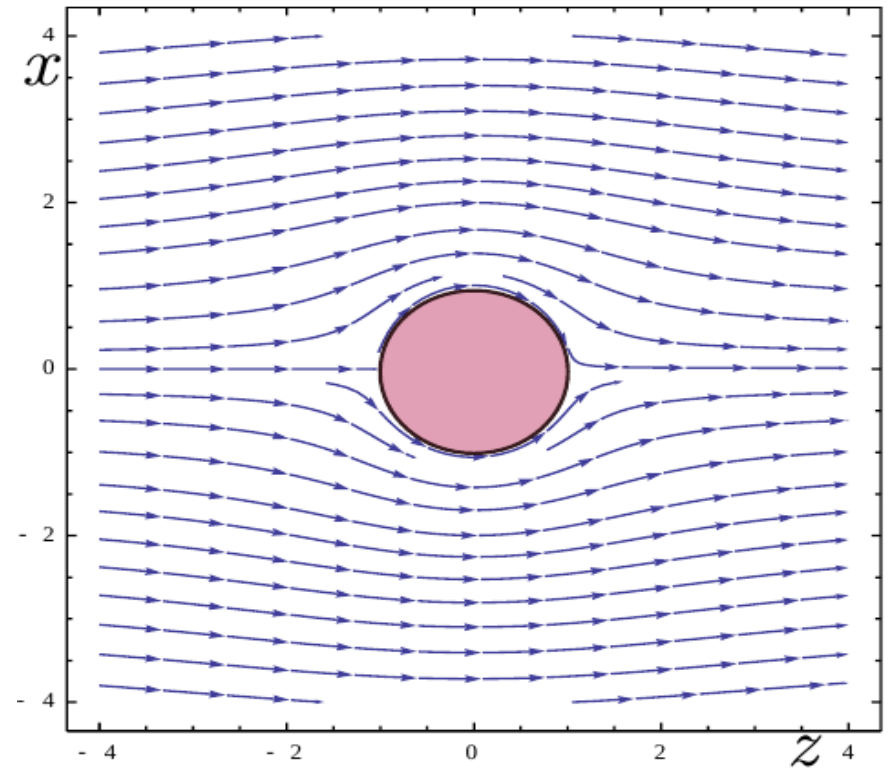
$$u_i(\mathbf{x}) = Q u_i^s + \sum_j P_j \frac{\partial}{\partial x_j} u_i^s + \sum_{j,k} T_{jk} \frac{\partial^2}{\partial x_j \partial x_k} u_i^s + \dots$$

$$\frac{\partial}{\partial x_j} u_i^s = \frac{1}{4\pi r^3} (\delta_{ij} - 3\hat{r}_i \hat{r}_j)$$

Moving Sphere:

$$u_i(\mathbf{x}) = \sum_j \left(b_j G_{ij} + P_j \frac{\partial}{\partial x_j} u_i^s \right)$$

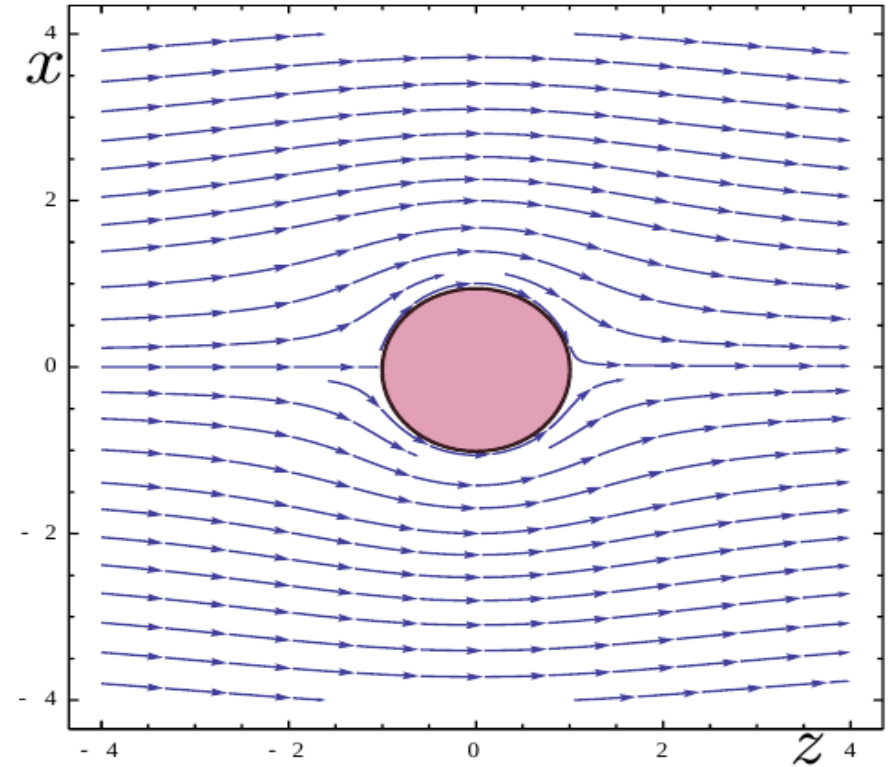
$$\mathbf{u}(\mathbf{x})|_{\mathbf{x}=R} = \mathbf{V}_0.$$



Moving Sphere:

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$$\mathbf{u}(\mathbf{x})|_{\mathbf{x}=R} = \mathbf{V}.$$



$$\mathbf{P} = \pi R^3 \mathbf{V}, \quad \mathbf{b} = \frac{4}{3} \pi \eta R \mathbf{V}.$$

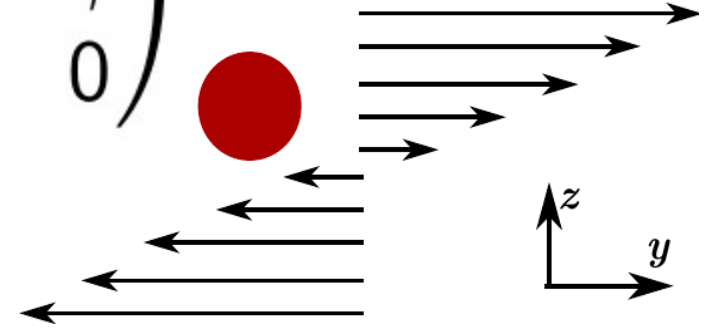
$$\mathbf{u}(\mathbf{r}) = \left(\frac{3}{4} \left(\frac{R}{r} \right) (I + \hat{r}\hat{r}) + \frac{1}{4} \left(\frac{R}{r} \right)^3 (I - 3\hat{r}\hat{r}) \right) \cdot \mathbf{V}.$$

Sphere in shear flow:

$$\mathbf{u}^\infty(\mathbf{x}) = \mathbf{\Gamma} \cdot \mathbf{x}$$

$$\mathbf{u}(\mathbf{x}) = \mathbf{u}^\infty(\mathbf{x}) + \mathbf{u}'(\mathbf{x})$$

$$\mathbf{\Gamma} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & \dot{\gamma} \\ 0 & 0 & 0 \end{pmatrix}$$

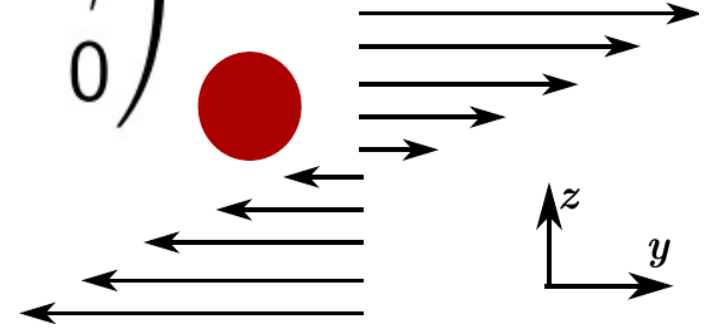


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$$\mathbf{u}'(\mathbf{x}) = 0, \quad x \rightarrow \infty$$

$$\mathbf{u}'(\mathbf{x}) = -\mathbf{\Gamma} \cdot \mathbf{x}, \quad x \rightarrow a$$

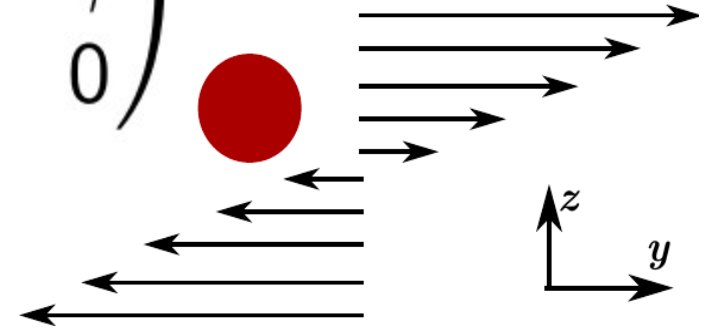
$$\mathbf{u}'(\mathbf{x}) = \text{Force Dipole} + \text{Source Quadrapole}$$

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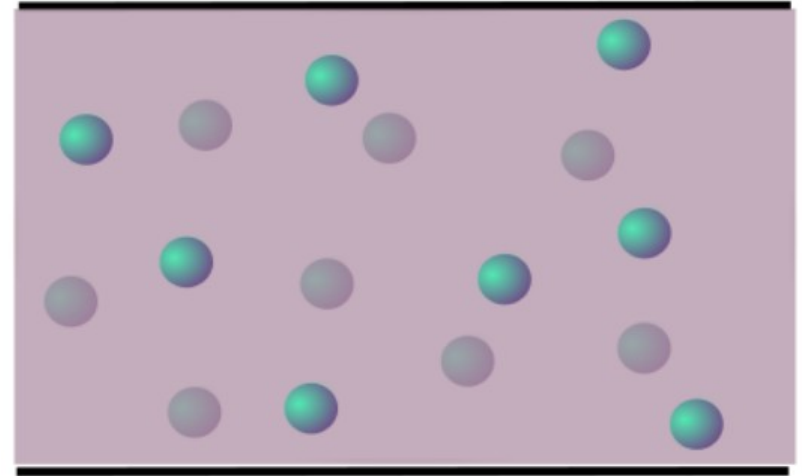
$$u'_i = G_{ijl}^D d_{jl} + Q_{ijl} q_{jl} \quad G_{ijl}^D = \partial_l G_{ij} \quad Q_{ijl} = \partial_j \partial_l u_i^s$$

$$\mathbf{d} = -\frac{5}{a^2} \mathbf{q} = -\frac{a^3}{6} (4\boldsymbol{\Gamma} + \boldsymbol{\Gamma}^T)$$

$$S = 8\pi\eta \frac{1}{2} (\mathbf{d} + \mathbf{d}^T) = -\frac{20}{6} \eta a^3 (\boldsymbol{\Gamma} + \boldsymbol{\Gamma}^T)$$


Passive suspension: Einstein's theory

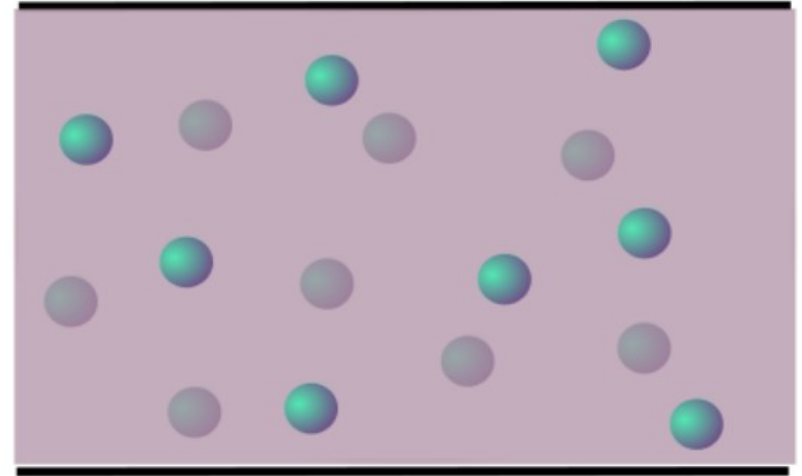
$$\eta^{\text{eff}} \sim \frac{\Sigma}{\nabla \mathbf{u}} \quad ?$$




In the presence of colloids, the applied force should do much work to establish the same velocity profile as in the bare fluid

Passive suspension: Einstein's theory

$$\eta^{\text{eff}} \sim \frac{\Sigma}{\nabla \mathbf{u}} \quad ?$$





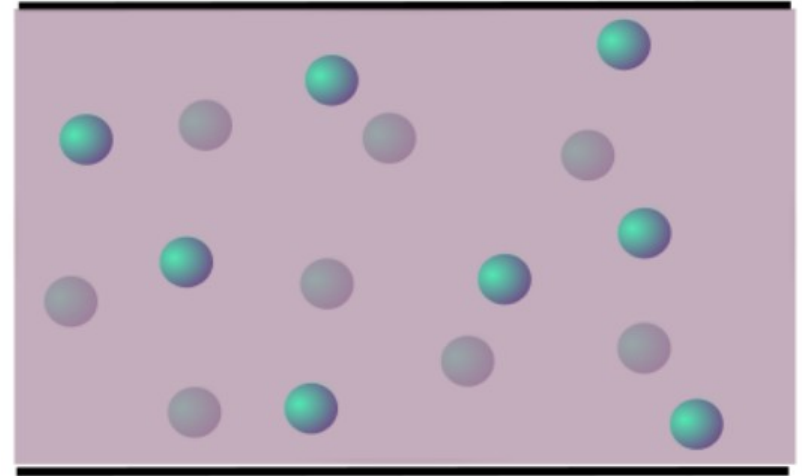
In the presence of colloids, the applied force should do much work to establish the same velocity profile as in the bare fluid

$$\eta^{\text{eff}} = \eta \left(1 + \frac{5}{2}c + \mathcal{O}(c)^2 \right)$$

Einstein 1906

Passive suspension: Einstein's theory

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In the presence of colloids, the applied force should do much work to establish the same velocity profile as in the bare fluid

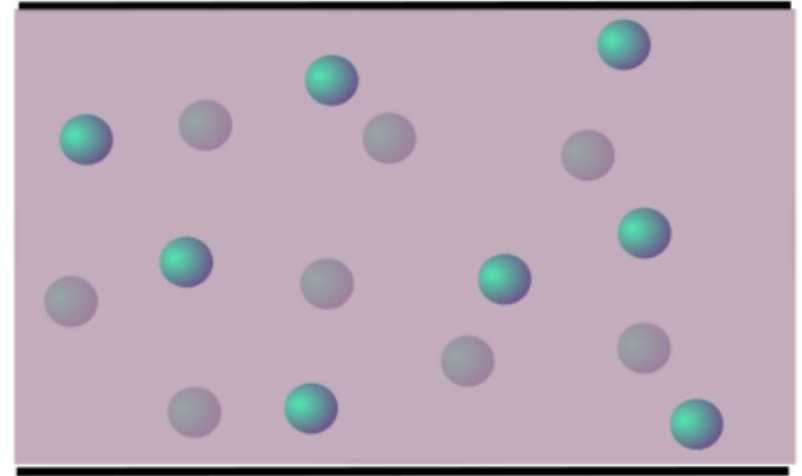
Perrin+Einstein 1909

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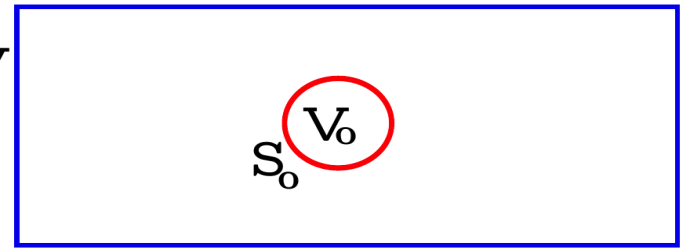
Einstein 1906

$$\eta^{\text{eff}} = \eta \left(1 + \frac{5}{2}c + \mathcal{O}(c)^2 \right)$$

*Interactions:
Batchelor 1972*

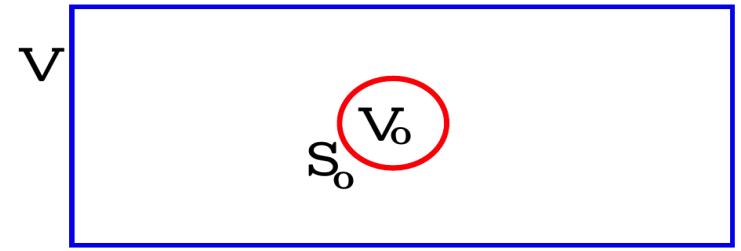
*Fluctuations:
Batchelor 1977*

$$\sigma_{ij} = -P\delta_{ij} + \eta \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

 V 

$$\langle \sigma_{ij} \rangle = \frac{1}{V} \int_V \sigma_{ij} dv = \frac{1}{V} \int_{V-V_0} \sigma_{ij} dv + \frac{1}{V} \int_{V_0} \sigma_{ij} dv$$

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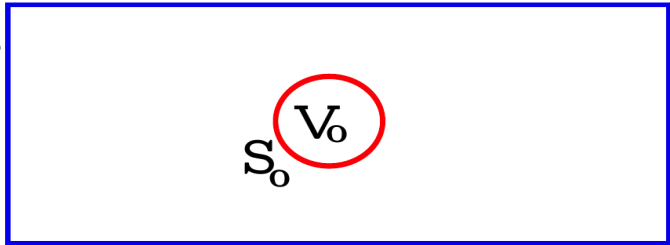
$$\int_V \partial_k (\sigma_{kj} x_i) = \int_V (\partial_k \sigma_{kj}) x_i + \int_V \sigma_{kj} \delta_{ki}$$

$$\oint_S \sigma_{kj} x_i n_k = 0 + \int_V \sigma_{ij}$$

$$= \frac{1}{V} \int_{V-V_0} \sigma_{ij} dv + \frac{1}{V} \int_{S_0} x_i \sigma_{jk} n_k ds$$

$$\sigma_{ij} = -P\delta_{ij} + \eta \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

V



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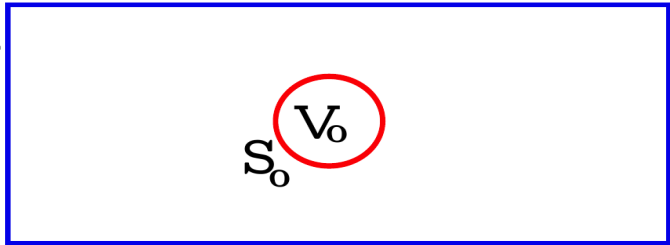
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$$= \frac{1}{V} \int_{V-V_0} \sigma_{ij} dv + \frac{1}{V} \int_{S_0} x_i \sigma_{jk} n_k ds$$

$$= \frac{1}{V} \int_{V-V_0} (-P\delta_{ij} + \eta(\partial_i u_j + \partial_j u_i)) dv + \frac{1}{V} \int_{S_0} x_i f_j ds \quad f_j = \sigma_{jk} n_k$$

$$\sigma_{ij} = -P\delta_{ij} + \eta \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

V



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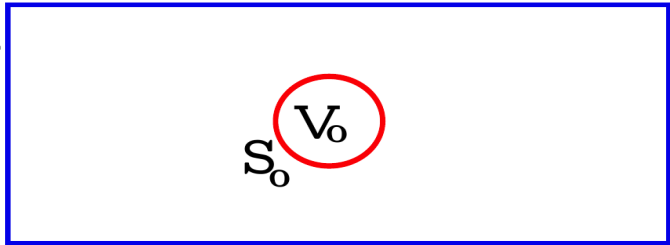
$$= \frac{1}{V} \int_{V-V_0} \sigma_{ij} dv + \frac{1}{V} \int_{S_0} x_i \sigma_{jk} n_k ds$$

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$$= \frac{1}{V} \delta_{ij} \int_{V-V_0} P dv + \eta \langle (\partial_i u_j + \partial_j u_i) \rangle + \frac{1}{V} \int_{S_0} (x_i f_j - \eta(u_i n_j + u_j n_i)) ds$$

$$\sigma_{ij} = -P\delta_{ij} + \eta \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

V



$$\langle \sigma_{ij} \rangle = \frac{1}{V} \int_V \sigma_{ij} dv = \frac{1}{V} \int_{V-V_0} \sigma_{ij} dv + \frac{1}{V} \int_{V_0} \sigma_{ij} dv$$

$$\int_V \partial_k (\sigma_{kj} x_i) = \int_V (\partial_k \sigma_{kj}) x_i + \int_V \sigma_{kj} \delta_{ki}$$

$$\oint_S \sigma_{kj} x_i n_k = 0 + \int_V \sigma_{ij}$$

$$= \frac{1}{V} \int_{V-V_0} \sigma_{ij} dv + \frac{1}{V} \int_{S_0} x_i \sigma_{jk} n_k ds$$

$$= \frac{1}{V} \int_{V-V_0} (-P\delta_{ij} + \eta(\partial_i u_j + \partial_j u_i)) dv + \frac{1}{V} \int_{S_0} x_i f_j ds \quad f_j = \sigma_{jk} n_k$$

$$= \frac{1}{V} \delta_{ij} \int_{V-V_0} P dv + \eta \langle (\partial_i u_j + \partial_j u_i) \rangle + \frac{1}{V} \int_{S_0} (x_i f_j - \eta(u_i n_j + u_j n_i)) ds$$

$$= \frac{1}{V} \delta_{ij} \int_{V-V_0} P dv + \eta \langle (\partial_i u_j + \partial_j u_i) \rangle + \frac{1}{V} S_{ij}$$

Force Dipole

Simple shear flow

$$\mathbf{u}(\mathbf{x}) = \mathbf{\Gamma} \cdot \mathbf{x}$$

$$\eta \langle (\partial_i u_j + \partial_j u_i) \rangle + \frac{1}{V} S_{ij} = \eta (\mathbf{\Gamma} + \mathbf{\Gamma}^T) + 0$$

Simple shear flow

$$\mathbf{u}(\mathbf{x}) = \boldsymbol{\Gamma} \cdot \mathbf{x}$$

$$\eta \langle (\partial_i u_j + \partial_j u_i) \rangle + \frac{1}{V} S_{ij} = \eta (\boldsymbol{\Gamma} + \boldsymbol{\Gamma}^T) + 0$$

Simple shear flow + Sphere

$$\eta \langle (\partial_i u_j + \partial_j u_i) \rangle + \frac{1}{V} S_{ij} = \eta (\boldsymbol{\Gamma} + \boldsymbol{\Gamma}^T) + \frac{1}{V} \frac{20\pi}{6} \eta a^3 (\boldsymbol{\Gamma} + \boldsymbol{\Gamma}^T)$$

$$\eta \rightarrow (1 + 5/2\phi)\eta$$

$$\phi = N \left(\frac{4}{3} \pi a^3 \right) / (V)$$

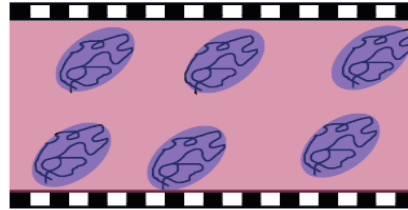
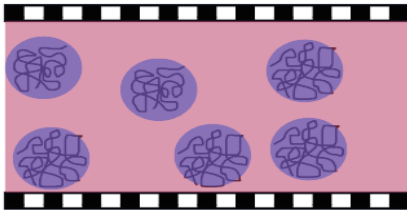
Non Newtonian Behavior:

Microscopic anisotropy

I: Intrinsic → Liquid Crystal



II: Induced anisotropy → Polymeric solutions



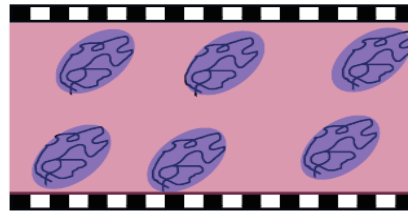
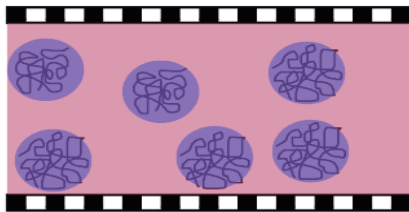
Non Newtonian Behavior:

Microscopic anisotropy

I: Intrinsic → Liquid Crystal



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Molecular relaxation time
Time scale for external force
 $\dot{\gamma}^{-1} < \tau \rightarrow \text{N-N}$

Molecules will not have enough time to
response and reach their isotropic structure

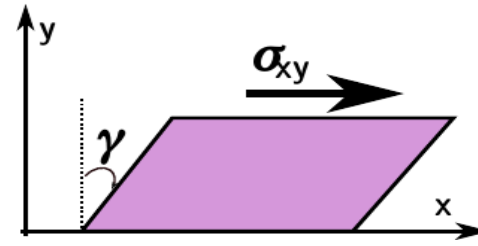
Relaxation time for water = 10^{-12} s

We expect to see N-N behavior in water for frequencies $> 10^{12}$ Hz

Phenomenological Description:

$$\sigma_{xy} = \eta(\dot{\gamma})\dot{\gamma}, \quad \sigma_{xz} = \sigma_{zy} = 0,$$

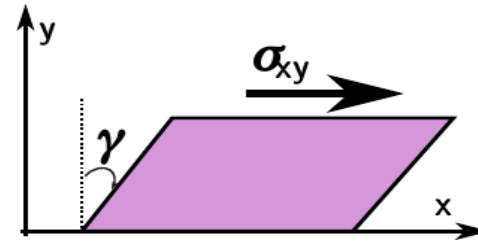
$$\sigma_{xx} - \sigma_{yy} = N_1(\dot{\gamma}), \quad \sigma_{yy} - \sigma_{zz} = N_2(\dot{\gamma})$$



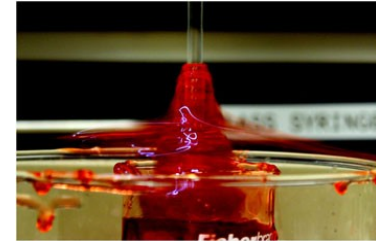
Phenomenological Description:

$$\sigma_{xy} = \eta(\dot{\gamma})\dot{\gamma}, \quad \sigma_{xz} = \sigma_{zy} = 0,$$

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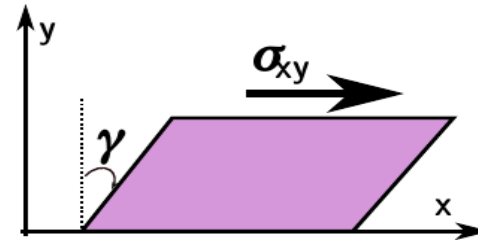
Normal Stress Differences Macroscopic Manifestation:
Weissenberg Effect



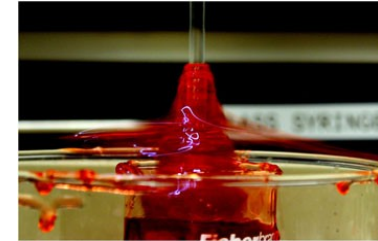
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Normal Stress Differences Macroscopic Manifestation: Weissenberg Effect



Shear thickening: Viscosity increases by increasing external shear
برش وشکسان: نشاسته ذرت و آب

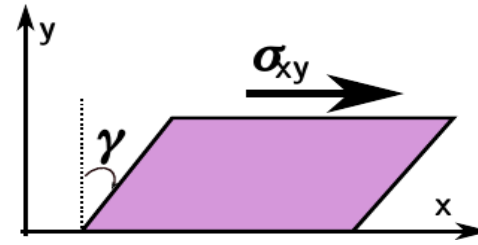
Shear thinning: Viscosity decreases by increasing external shear
برش روان: سس گوجه



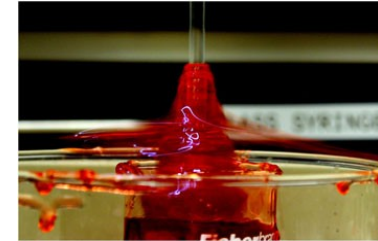
Phenomenological Description:

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Normal Stress Differences Macroscopic Manifestation: Weissenberg Effect



Shear thickening: Viscosity increases by increasing external shear
برش وشکسان: نشاسته ذرت و آب

Shear thinning: Viscosity decreases by increasing external shear
برش روان: سس گوجه



Viscoelastic Behavior → Kelvin Voigt Model

$$\sigma = G\gamma + \eta\dot{\gamma}$$

$$\gamma(t) = \gamma_0 e^{-i\omega t}$$

$$\sigma(t) = \tilde{G}(\omega)\gamma(t)$$

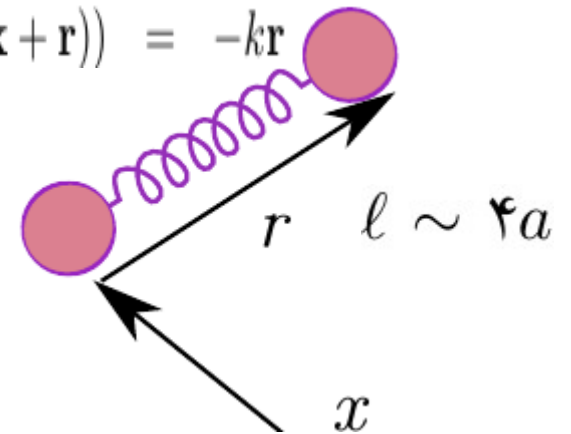
$$\tilde{G}(\omega) = G - i\omega\eta$$

Molecular Anisotropy: Microscopic Model

$$\zeta\pi\eta a(\mathbf{V}_\parallel - \mathbf{u}(\mathbf{x})) = k\mathbf{r},$$

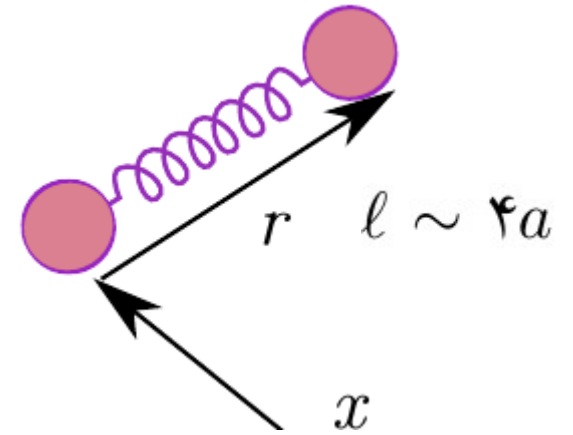
$$\zeta\pi\eta a(\mathbf{V}_\parallel - \mathbf{u}(\mathbf{x} + \mathbf{r})) = -k\mathbf{r}$$

$$\zeta\pi\eta a(\mathbf{V}_\parallel - \mathbf{V}_\parallel - \mathbf{u}(\mathbf{x} + \mathbf{r}) + \mathbf{u}(\mathbf{x})) = -\zeta k\mathbf{r}$$



Molecular Anisotropy: Microscopic Model

$$\begin{aligned} \xi\pi\eta a(\mathbf{V}_\gamma - \mathbf{V}_\gamma - \mathbf{u}(x + \mathbf{r}) + \mathbf{u}(x)) &= -\gamma k\mathbf{r} \\ \xi\pi\eta a \frac{d}{dt}\mathbf{r} &= -\gamma k\mathbf{r} + \xi\pi\eta a(\mathbf{r} \cdot \nabla)\mathbf{u}(\mathbf{x}) \end{aligned}$$



Molecular Anisotropy: Microscopic Model

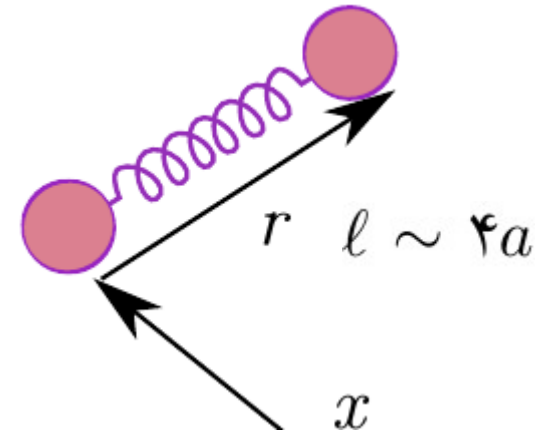
$$\zeta\pi\eta a(\mathbf{V}_\zeta - \mathbf{V}_\zeta - \mathbf{u}(x + \mathbf{r}) + \mathbf{u}(x)) = -\zeta k\mathbf{r}$$

$$\zeta\pi\eta a \frac{d}{dt} \mathbf{r} = -\zeta k\mathbf{r} + \zeta\pi\eta a(\mathbf{r} \cdot \nabla)\mathbf{u}(\mathbf{x})$$

$$\frac{d}{dt} \mathbf{r} = -\frac{1}{\zeta\tau} \mathbf{r} + (\mathbf{r} \cdot \nabla)\mathbf{u}(\mathbf{x}) + \left(\frac{1}{\zeta\pi\eta a}\right)\mathbf{f}(t)$$

$$\langle f_i(t)f_j(t') \rangle = \zeta\xi k_B T \delta_{ij} \delta(t - t')$$

$$\tau = \zeta\pi\eta a / \zeta k$$



Molecular Anisotropy: Microscopic Model

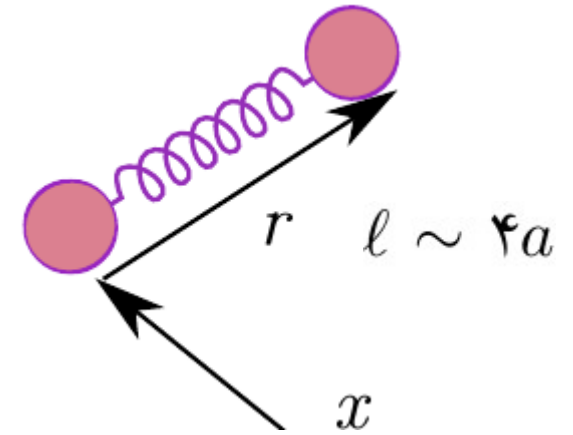
$$\zeta\pi\eta a(\mathbf{V}_\zeta - \mathbf{V}_\zeta - \mathbf{u}(x + \mathbf{r}) + \mathbf{u}(x)) = -\zeta k\mathbf{r}$$

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$$\sigma^P \cdot d\mathbf{s} = (k\mathbf{r})(\mathcal{N}\mathbf{r} \cdot d\mathbf{s}) \quad \sigma^P = \mathcal{N}k\mathbf{r}\mathbf{r} \quad \mathbf{r} = \mathbf{r}(\mathbf{x}, t)$$

Molecular Anisotropy: Microscopic Model

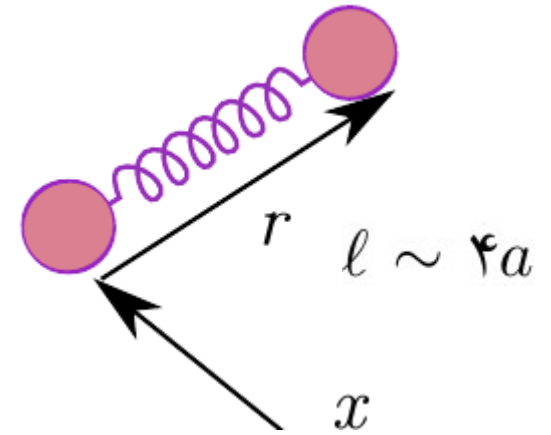
$$\zeta\pi\eta a(\mathbf{V}_\zeta - \mathbf{V}_\zeta - \mathbf{u}(x + \mathbf{r}) + \mathbf{u}(x)) = -\zeta k\mathbf{r}$$

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Molecular Anisotropy: Microscopic Model

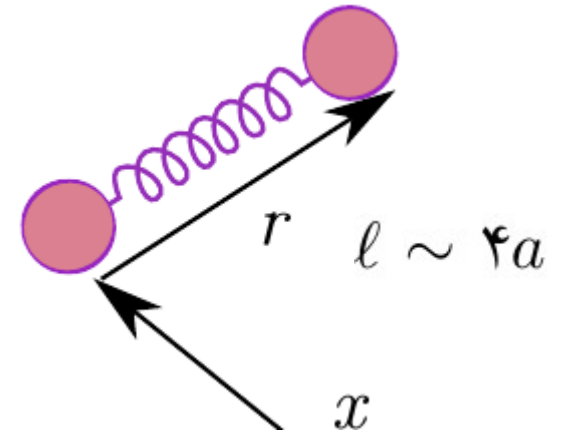
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$$\langle f_i(t)f_j(t') \rangle = \zeta\xi k_B T \delta_{ij} \delta(t - t')$$

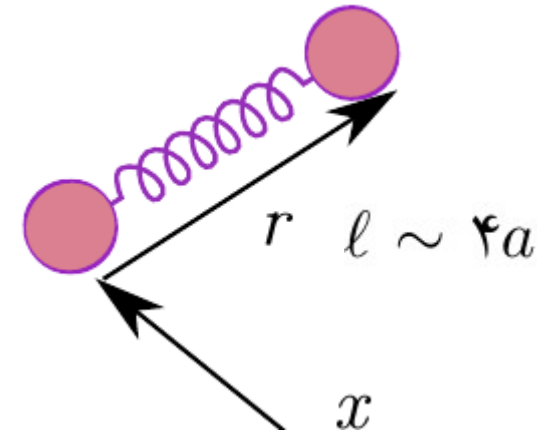
$$\tau = \zeta\pi\eta a / \zeta k$$



$$\sigma^P \cdot d\mathbf{s} = (k\mathbf{r})(\mathcal{N}\mathbf{r} \cdot d\mathbf{s}) \quad \sigma^P = \mathcal{N}k\mathbf{r}\mathbf{r} \quad \mathbf{r} = \mathbf{r}(\mathbf{x}, t)$$

Molecular Anisotropy: Microscopic Model

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$$\begin{aligned} \frac{d}{dt} \mathbf{r} &= -\frac{1}{\zeta \tau} \mathbf{r} + (\mathbf{r} \cdot \nabla) \mathbf{u}(\mathbf{x}) + \left(\frac{1}{\zeta \pi \eta a} \right) \mathbf{f}(t) \\ \langle f_i(t) f_j(t') \rangle &= \zeta \xi k_B T \delta_{ij} \delta(t - t') \end{aligned} \quad \tau = \zeta \pi \eta a / \zeta k$$

$$\sigma^P \cdot d\mathbf{s} = (k\mathbf{r})(\mathcal{N}\mathbf{r} \cdot d\mathbf{s}) \quad \sigma^P = \mathcal{N}k\mathbf{r}\mathbf{r} \quad \mathbf{r} = \mathbf{r}(\mathbf{x}, t)$$

After averaging over noise, it is easy to show: $\Sigma = \frac{\zeta}{\mathcal{N}k_B T} \sigma^P$

$$\left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) \Sigma - \Sigma \cdot \nabla \mathbf{u} - \nabla \mathbf{u}^T \cdot \Sigma = -\frac{1}{\tau} (\Sigma - I)$$

Oldroyd Model

$$\left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) \mathbf{u} = \nabla \cdot \sigma, \quad \nabla \cdot \mathbf{u} = 0,$$

$$\left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) \Sigma - \Sigma \cdot \nabla \mathbf{u} - \nabla \mathbf{u}^T \cdot \Sigma = -\frac{1}{\tau} (\Sigma - I)$$

$$\sigma = -PI + \eta(\nabla \mathbf{u} + \nabla \mathbf{u}^T) + g\Sigma \quad g = \frac{1}{\tau} \mathcal{N} k_B T$$

Oldroyd Model

$$\left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) \mathbf{u} = \nabla \cdot \sigma, \quad \nabla \cdot \mathbf{u} = 0,$$

$$\left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) \Sigma - \Sigma \cdot \nabla \mathbf{u} - \nabla \mathbf{u}^T \cdot \Sigma = -\frac{1}{\tau} (\Sigma - I)$$

$$\sigma = -PI + \eta(\nabla \mathbf{u} + \nabla \mathbf{u}^T) + g\Sigma \quad g = \frac{1}{\tau} \mathcal{N} k_B T$$

Response to a simple shear flow:

$$\mathbf{u} = (\dot{\gamma}y, 0, 0) \quad \nabla \mathbf{u} = \begin{pmatrix} 0 & 0 \\ \dot{\gamma} & 0 \end{pmatrix} \quad \Sigma = \begin{pmatrix} \Sigma_{xx} & \Sigma_{xy} \\ \Sigma_{xy} & \Sigma_{yy} \end{pmatrix}$$

Oldroyd Model

$$\left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) \mathbf{u} = \nabla \cdot \sigma, \quad \nabla \cdot \mathbf{u} = 0,$$

$$\left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) \Sigma - \Sigma \cdot \nabla \mathbf{u} - \nabla \mathbf{u}^T \cdot \Sigma = -\frac{1}{\tau} (\Sigma - I)$$

$$\sigma = -PI + \eta(\nabla \mathbf{u} + \nabla \mathbf{u}^T) + g\Sigma \quad g = \frac{1}{\tau} \mathcal{N} k_B T$$

Response to a simple shear flow:

$$\mathbf{u} = (\dot{\gamma}y, 0, 0) \quad \nabla \mathbf{u} = \begin{pmatrix} 0 & 0 \\ \dot{\gamma} & 0 \end{pmatrix} \quad \Sigma = \begin{pmatrix} \Sigma_{xx} & \Sigma_{xy} \\ \Sigma_{xy} & \Sigma_{yy} \end{pmatrix}$$

$$\dot{\gamma} \begin{pmatrix} 2\Sigma_{xy} & \Sigma_{yy} \\ \Sigma_{yy} & 0 \end{pmatrix} = \frac{1}{\tau} \begin{pmatrix} \Sigma_{xx} - 1 & \Sigma_{xy} \\ \Sigma_{xy} & \Sigma_{yy} - 1 \end{pmatrix}$$

Oldroyd Model

$$\left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) \mathbf{u} = \nabla \cdot \sigma, \quad \nabla \cdot \mathbf{u} = 0,$$

$$\left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) \Sigma - \Sigma \cdot \nabla \mathbf{u} - \nabla \mathbf{u}^T \cdot \Sigma = -\frac{\lambda}{\tau} (\Sigma - I)$$

$$\sigma = -PI + \eta(\nabla \mathbf{u} + \nabla \mathbf{u}^T) + g\Sigma \quad g = \frac{\lambda}{\tau} \mathcal{N} k_B T$$

Response to a simple shear flow:

$$\mathbf{u} = (\dot{\gamma}y, 0, 0) \quad \nabla \mathbf{u} = \begin{pmatrix} 0 & \dot{\gamma} \\ \dot{\gamma} & 0 \end{pmatrix} \quad \Sigma = \begin{pmatrix} \Sigma_{xx} & \Sigma_{xy} \\ \Sigma_{xy} & \Sigma_{yy} \end{pmatrix}$$

$$\dot{\gamma} \begin{pmatrix} \Sigma_{xy} & \Sigma_{yy} \\ \Sigma_{xy} & 0 \end{pmatrix} = \frac{\lambda}{\tau} \begin{pmatrix} \Sigma_{xx} - 1 & \Sigma_{xy} \\ \Sigma_{xy} & \Sigma_{yy} - 1 \end{pmatrix}$$

$$\Sigma_{xx} = 1 + 2\dot{\gamma}^2 \tau^2, \quad \Sigma_{xy} = \dot{\gamma} \tau, \quad \Sigma_{yy} = 1 \quad \tau = 9\pi\eta a / 4k$$

$$\sigma = \begin{pmatrix} -P & \eta\dot{\gamma} \\ \eta\dot{\gamma} & -P \end{pmatrix} + g \begin{pmatrix} 1 + 2\dot{\gamma}^2 \tau^2 & \dot{\gamma} \tau \\ \dot{\gamma} \tau & 1 \end{pmatrix}$$

$$N_1 = \sigma_{xx} - \sigma_{yy} = 2g\dot{\gamma}^2 \tau^2$$

$$\eta_e = \eta + g\tau = \eta(1 + \epsilon)$$

$$\epsilon = (3/4)\pi \mathcal{N} k_B T a / k$$

Response to an harmonic shear flow:

$$\mathbf{u} = (\dot{\gamma}y, 0, 0)$$

$$\dot{\gamma}(t) = a\omega \cos \omega t$$

$$\sigma_{xy} = \tilde{G}(\omega)\gamma(t)$$

$$\tilde{G}(\omega) = G_r(\omega) + iG_i(\omega)$$

$$G_r(\omega) = ?$$

$$G_i(\omega) = ?$$

از توجه شما ممنونم